

# TANTALUM CAPACITOR AS A MIS STRUCTURE: TRANSPORT CHARACTERISTICS TEMPERATURE DEPENDENCIES

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**Abstract:** Temperature dependencies of a leakage current in normal mode are explained on the basis of a model, in which the solid state tantalum capacitor is considered as a metal-insulator-semiconductor (MIS) heterostructure. The measurement was performed in temperature range from 105 to 155°C. Ohmic conductivity increases exponentially with increasing temperature with activation energy 0.94 eV. Tunneling voltage parameter and tunneling energy barrier decreases with increasing temperature, reaching values 0.45 to 0.26 eV.

**Keywords:** tantalum capacitor, MIS structure,  $I$ - $V$  characteristics, tunneling energy barrier, activation energy

## 1. INTRODUCTION

A tantalum capacitor consists of a metallic Ta anode, insulator-Ta<sub>2</sub>O<sub>5</sub> film and a semiconductor-MnO<sub>2</sub> cathode, which is made of carbon and dipping silver layers. The ideal metal-insulator-semiconductor (MIS) structure theory [1], to be considered in this paper, serves as a foundation for understanding the real tantalum capacitor characteristics. Current through an amorphous insulating film increases roughly exponentially with applied voltage [2].

Mechanisms, which may explain the observed characteristics, are: Ohmic electron transport, Poole-Frenkel electron transport with distributed traps, and tunneling electron transport.

Leakage current in our samples may be described as an electron transport with Ohmic, Poole-Frenkel, and tunneling components.

$I$ - $V$  characteristics have three components given by:

$$I = G_{\Omega}V + G_{PF}V \exp(\beta_{PF} \sqrt{V}) + G_T V^2 \exp(-V/V_T), \quad (1)$$

where  $G_{\Omega}$  is Ohmic conductivity,  $G_{PF}$  is Poole-Frenkel conductivity,  $\beta_{PF}$  is Poole-Frenkel coefficient,  $G_T$  is tunneling transparencies or tunneling conductivity per volt in S/V, and  $V_T$  is tunneling voltage parameter.

Ohmic conductivity  $G_{\Omega}$  is given by:

$$G_{\Omega} = Ae \mu n / d, \quad (2)$$

where  $A$  is the area of capacitor,  $e$  is its elementary charge,  $\mu$  is electron mobility,  $n$  is electron concentration and  $d$  is insulating layers thickness.

Electron concentration is exponential function of reciprocal given by:

$$n = n_0 \exp(-E_a/kT), \quad (3)$$

where  $n_0$  is electron concentration at temperature  $T_0$ ,  $E_a$  is activation energy and  $kT$  is electron thermal energy. Electron mobility and other quantities can be in short temperature range considered to be constants.

Then we can write that:

$$G_\Omega = G_{\Omega 0} \exp(-E_a/kT), \quad (4)$$

where

$$G_{\Omega 0} = Ae \mu_{n0} / d. \quad (5)$$

Activation energy  $E_a$  can be obtained from temperature dependence of Ohmic conductivity.

Poole-Frenkel coefficient  $\beta_{PF}$  is given by

$$\beta_{PF} = 2(e^3 / \pi \epsilon_0 \epsilon_r d)^{1/2} / kT, \quad (6)$$

where  $e$  is elementary charge,  $\epsilon_0$  is permittivity of vacuum,  $\epsilon_r$  is Ta<sub>2</sub>O<sub>5</sub> layer permittivity.

Poole-Frenkel coefficient vs. temperature for insulating layer thickness  $d = 210$  nm is given by:

$$\beta_{PF} = 362.12 / T, \quad (7)$$

where  $T$  is temperature in Kelvin.

**Table 1:** Poole-Frenkel coefficient vs. temperature

T / °C	23	85	105	125	135	145	155
T / K	296.15	358.15	378.15	398.15	408.15	418.15	428.15
$\beta_{PF} / V^{-1/2}$	1.21	1.0	0.96	0.9	0.88	0.86	0.84

### The impurity states in the forbidden energy gap

Insulating Ta<sub>2</sub>O<sub>5</sub> films always contain impurities like oxygen vacancies and traps. These impurities form localized states for charge carriers in the forbidden energy gap. When the localized electronic wave functions of such impurities overlap, an electron bound to one impurity state can tunnel to an unoccupied state without involving activation into the conduction band. This tunneling process between impurity sites is referred to as impurity conduction.

### Electrical transport by hopping and tunneling

The mobility of an electron moving through impurity states is small especially at low temperatures. Then, the conduction mechanism is mainly influenced by electron hopping between neighboring impurity site. This type of conduction process depends on impurity concentration and the energy depths of impurity states. The concept of hopping transport has been used for a long time in connection with ionic conduction; ions move essentially by hopping, whether through interstices or vacancies.

The leakage current is at high electric fields mainly influenced by electron tunneling through a potential barrier between Ta<sub>2</sub>O<sub>5</sub> insulating layer and cathode or anode. The potential barriers can be approximated by a square or triangular shapes (see Figure 1). The triangular shape facilitates an inter-molecular electron transfer, since the barrier width at higher fields becomes effectively smaller for high energy electrons.

## Tunneling current

The charge carrier transport in tantalum pentoxide thin film at high electric fields becomes mainly given by the tunneling current [3 and 4]. The barrier height depends on the work function of the cathode material and the electron affinity of the Ta<sub>2</sub>O<sub>5</sub> insulating layer, as it is shown in Figure 1 for a triangular potential barrier.

In normal mode (for Ta electrode positive) the  $I$ - $V$  characteristic for tunneling current is described by the relation:

$$I = G_T V^n \exp\left(-\frac{V_T}{V}\right) \quad (8)$$

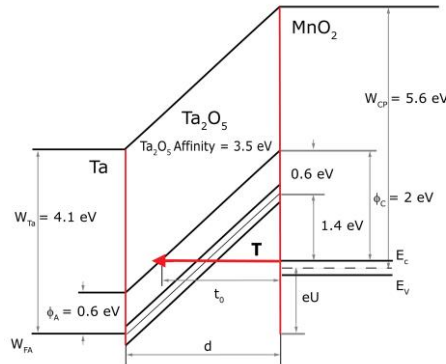
where  $V$  is the applied voltage,  $G_T$  and  $V_T$  are the constants, and the value of exponent  $n$  depends on the barrier shape (for triangular barrier  $n = 2$ ).

The tunneling parameter  $V_T$  can be expressed for the energy barrier  $e\Phi_T$  and the thickness  $t_0$  of the insulating layer as

$$V_T = \left( \frac{8\pi\sqrt{2m^*}}{(3eh)(e\Phi_0)^{1.5}t_0} \right), \quad (9)$$

here,  $m^*$  is the electron effective mass,  $h = 6.6 \times 10^{-34}$  Js is the Planck constant.

Evaluating the potential barrier from  $I$ - $V$  characteristic we can further determine the parameter  $V_T$  for the potential barrier height  $\Phi_0$ .



**Figure 1:** The triangular potential barrier model for Ta capacitor with MnO<sub>2</sub> cathode

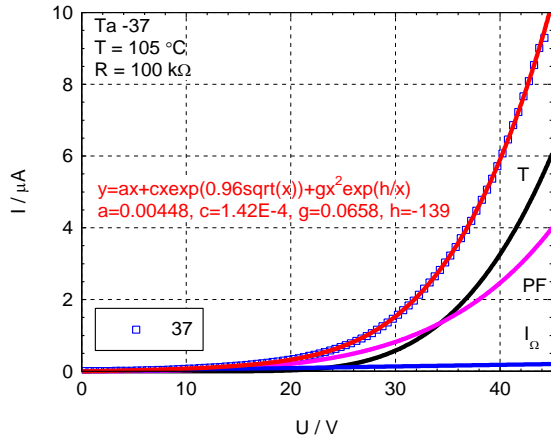
## 2. EXPERIMENTAL

To get more information on the current transport processes, the  $I$ - $V$  characteristics were measured in normal mode (Ta electrode is positive) in the temperature range from  $T = 105$  to  $155$  °C. The fitting of experimental dependencies were performed in order to find the parameters important for the construction of the energy diagram of the MIS structure. Studied capacitors have Ta<sub>2</sub>O<sub>5</sub> layer prepared by anodic oxidation. Measured sample TaS4 – 37 has oxide thickness  $d = 210$  nm, AC capacitance  $C = 10$  uF/25V, and rated voltage 35 V.

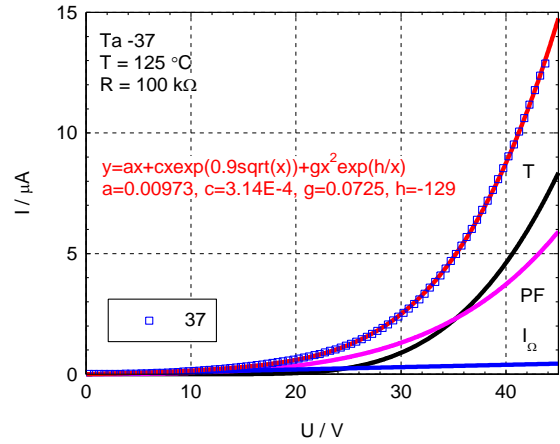
## 3. DISCUSSION

$I$ - $V$  characteristics of sample TaS4 - 37 measured at temperatures  $T = 105$  to  $155$  °C are shown in Figures 2 to 5.  $I$ - $V$  characteristics were measured with voltage step 0.1 V and time delay between two successive steps 20 s to have low value of polarization current component, which was removed from measured data. Analysis was performed by fitting of experimental data using equation (1). This equation contains 6 parameters. Therefore the best coincidence of fitting curve with

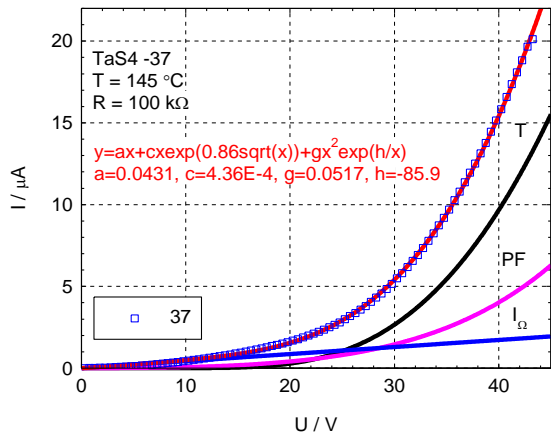
experimental data in all voltage ranges is difficult. For this reason we have used theoretical value of Poole-Frenkel coefficient  $\beta_{PF}$  to decrease the number parameters. It is due to that Poole-Frenkel current component is at rated voltage lower than tunneling one at all measured temperature ranges.



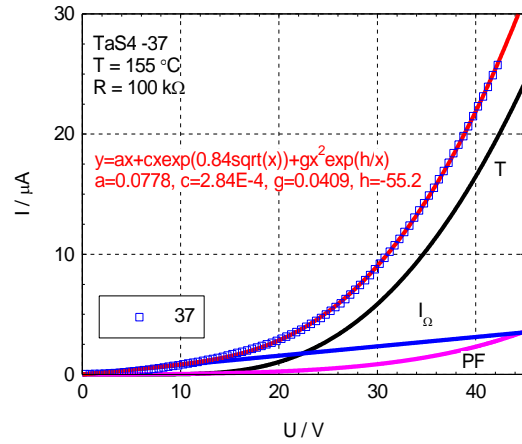
**Figure 2:** *I-V* characteristics of sample TaS4 - 37 measured at temperature  $T = 105\text{ °C}$



**Figure 3:** *I-V* characteristics of sample TaS4 - 37 measured at temperature  $T = 125\text{ °C}$



**Figure 4:** *I-V* characteristics of sample TaS4 - 37 measured at temperature  $T = 145\text{ °C}$

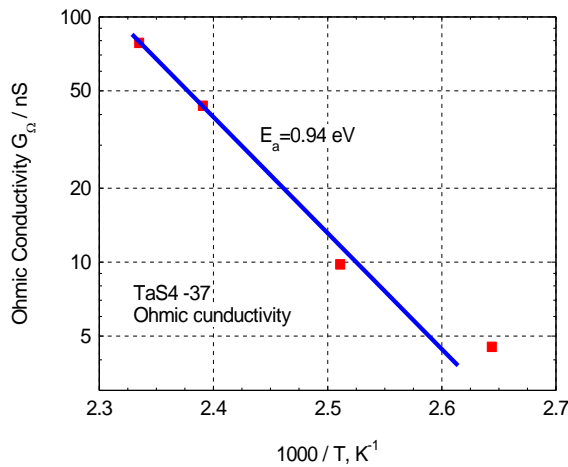


**Figure 5:** *I-V* characteristics of sample TaS4 - 37 measured at temperature  $T = 155\text{ °C}$

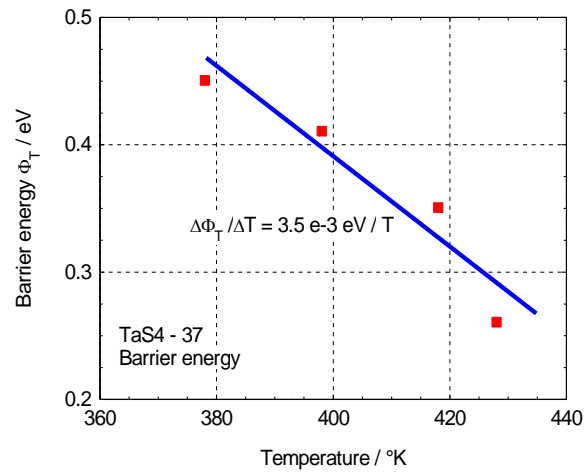
**Table 2:** *I-V* characteristics parameters of sample TaS4-37

Temperature: $T / \text{°C}$	105	125	145	155
Absolute temperature: $T / \text{K}$	378.15	398.15	418.15	428.15
Ohmic conductivity: $G_{\Omega} / \text{nS}$	4.48	9.73	43.1	77.8
P-F conductivity: $G_{PF} / \text{pS}$	142	314	436	284
P-F coefficient: $G_{PF} / V^{-0.5}$	0.96	0.9	0.86	0.84
Tunnel transparencies: $G_T / \text{nS/V}$	65.8	72.5	51.7	40.9
Tunneling voltage: $V_T / \text{V}$	139	129	85.9	55.2
Barrier energy: $\Phi_T / \text{eV}$	0.45	0.41	0.35	0.26

We will give graphical evolution of Ohmic conductivity vs. reciprocal temperature (see Figure 6) and in Figure 7 tunneling barrier energy  $\Phi_T$  vs. temperature.



**Figure 6:** Activation energy of Ohmic conductivity for sample TaS4 - 37



**Figure 7:** Tunneling barrier energy  $\Phi_T$  vs. temperature for sample TaS4 - 37

#### 4. CONCLUSION

Ohmic conductivity vs. reciprocal temperature give an activation energy  $E_a = 0.94$  eV which is equal to the energy of oxygen vacancies band below conduction band of Ta<sub>2</sub>O<sub>5</sub> insulating layer. This value is very near to that of other experiments (see [3 and 4]).

Tunneling voltage parameter and tunneling barrier energy decrease with increasing temperature. It means that with increasing temperature tunneling current components increases while Poole-Frenkel current component decreases.

Poole-Frenkel conductivity increases with increasing temperature. Some new effect at temperature 155 °C where Poole-Frenkel conductivity is lower than at temperature 125 °C is described. Similar effect is obtained for the tunneling barrier energy  $\Phi_T$  at temperature 155 °C. It can be caused by some change in oxygen vacancies distribution. Single  $I$ - $V$  characteristics measurement was performed during 2.5 h, then at high temperature ionic motion can affect value of current at high voltage range.

#### ACKNOWLEDGEMENT

This work was supported by the Internal Grant Agency of Brno University of Technology, grant No. FEKT-S-17-4626

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